Probability Operator Measure and Phase Measurement in a Deformed Hilbert Space

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We discuss the probability operator measure and phase measurement in a deformed Hilbert space.

1. INTRODUCTION

We consider the set

$$H_q = \{f: f(z) = \sum a_n z^n \text{ where } \sum [n]! |a_n|^2 < \infty\}$$

where $[n] = (1 - q^n)/(1 - q), 0 < q < 1.$

For $f, g \in H_q, f(z) = \sum_{n=0}^{\infty} a_n z^n$, and $g(z) = \sum_{n=0}^{\infty} b_n z^n$ we define addition and scalar multiplication as follows:

$$f(z) + g(z) = \sum_{n=0}^{\infty} (a_n + b_n) z^n$$
(1)

and

$$\lambda \circ f(z) = \sum_{n=0}^{\infty} \lambda a_n z^n$$
(2)

It is easily seen that H_q forms a vector space with respect to usual pointwise scalar multiplication and pointwise addition by (1) and (2). We observe that $e_q(z) = \sum_{n=0}^{\infty} (z^n/[n]!)$ belongs to H_q .

Now we define the inner product of two functions $f(z) = \sum a_n z^n$ and $g(z) = \sum b_n z^n$ belonging to H_q as

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$$(f,g) = \sum [n]! \,\overline{a}_n b_n \tag{3}$$

The corresponding norm is given by

$$||f||^2 = (f, f) = \sum [n]! |a_n|^2 < \infty$$

With this norm derived from the inner product it can be shown that H_q is a complete normed space. Hence H_q forms a Hilbert space.

In a recent paper [1] we have proved that the set $\{f_n \equiv z^n/\sqrt{[n]!}, n = 0, 1, 2, 3, ...\}$ forms a complete orthonormal set. If we consider the following actions on H_q :

$$Tf_{n} = \sqrt{[n]} f_{n-1}$$

$$T^{*}f_{n} = \sqrt{[n+1]} f_{n+1}$$
(4)

where T is the backward shift and its adjoint T^* is the forward shift operator on H_q , then we have shown [1] that the solution of the eigenvalue equation

$$Tf = \alpha f \tag{5}$$

is given by

$$f_{\alpha} = e_q(|\alpha|^2)^{-1/2} \sum_{n=0}^{\infty} \frac{\alpha^n}{\sqrt{[n]!}} f_n$$
 (6)

We call f_{α} a coherent vector in H_q .

The paper is divided into five sections. In this section we have given an introduction stating coherent vectors in H_q . In Section 2 we describe the probability operator measure (POM) in H_q . In Section 3 we discuss phase distribution in H_q . In Section 4 we study the phase estimation problem, and in Section 5 we give a conclusion.

2. PROBABILITY OPERATOR MEASURE

A discrete spectrum *probability operator measure (POM)* consists of a set of Hermitian, positive-semidefinite operators $\{\Pi_n : n \in \mathbb{N}\}$ which resolves the identity

$$I = \sum_{n \in \mathbb{N}} \prod_{n} \tag{7}$$

Measurement of this POM, by definition, gives a discrete, classical random variable with probability distribution

$$P(n, g) = (g, \Pi_n g) \quad \text{for} \quad n \in \mathbb{N}$$
(8)

where g is any vector of unit norm in H_q .

In order that the laws of classical probability be satisfied, it is necessary and sufficient that

$$0 \le P(n, g) \le 1, \qquad \sum_{n=0}^{\infty} P(n, g) = 1$$
 (9)

are satisfied for arbitrary g of unit norm in H_q .

We know that the sequence $f_n = z^n / \sqrt{[n]!}$ forms a complete orthonormal sequence in H_q and comprises eigenvectors of the operator $N = T^*T$ such that

$$Nf_n = [n]f_n \tag{10}$$

Measurement of N for any arbitrary vector $g \in H_q$ of unit norm yields a discrete-valued classical random variable with probability distribution

$$P(f_n, g) = |(f_n, g)|^2$$
 for $n = 0, 1, 2, ...$ (11)

In order that the law of classical probability be satisfied, it is necessary and sufficient that

$$0 \le P(f_n, g) \le 1, \qquad \sum_{n=0}^{\infty} P(f_n, g) = 1$$
 (12)

for arbitrary $g \in H_q$ of unit norm.

The completeness of $\{f_n\}$ guarantees that the prescription in equation (11) obeys equation (12). For, if we expand the arbitrary vector g of unit norm in terms of f_n , we have

$$g = \sum_{n=0}^{\infty} (f_n, g) f_n$$
(13)
$$= \sum_{n=0}^{\infty} |f_n\rangle \langle f_n|g$$

Where we define the operator

$$|f_n\rangle\langle f_n|: H_q \to H_q$$

by

$$|f_n\rangle\langle f_n| = (f_n, g)f_n$$

Equation (12) is now easily verified from equation (11) and equation (13). Thus, N operator measurement is equivalent to the POM

$$\{\Pi_n = |f_n\rangle\langle f_n|: n = 0, 1, 2, \ldots\}$$
(14)

Similarly, a continuous spectrum POM consists of a set of Hermitian,

positive-semidefinite differential operators $\{d \Pi(\beta): \beta \in \mathbb{C}\}$ which resolve the identity

$$I = \int_{\beta \in \mathbb{C}} d\Pi(\beta) \tag{15}$$

The result of measuring this POM is, by definition, a continuous classical random variable whose probability density function is given by

$$p(\beta, g) = \frac{(g, d \Pi(\beta)g)}{d\beta} \quad \text{for} \quad \beta \in \mathbb{C}$$
 (16)

where g is any vector of unit norm in H_q .

We know that the backwardshift *T* has eigenvectors—the coherent vectors f_{α} (6). These vectors are not orthogonal, but they form a resolution of the identity

$$I = \frac{1}{2\pi} \int_{\alpha \in \mathbb{C}} d\mu(\alpha) |f_{\alpha}\rangle \langle f_{\alpha}|$$
(17)

where

$$d\mu(\alpha) = e_q(|\alpha|^2)e_q(-|\alpha|^2)d_q|\alpha|^2 \ d\theta$$

with $\alpha = re^{i\theta}$, which defines a *T*-POM

$$d\Pi(\alpha) \equiv |f_{\alpha}\rangle \langle f_{\alpha}| \frac{d\mu(\alpha)}{2\pi} \quad \text{for } \alpha \in \mathbb{C}$$
 (18)

The outcome of the *T*-POM is a complex-valued continuous classical random variable with probability density function

$$p(\alpha, g) = \frac{(g, d \Pi(\alpha)g)}{d\mu(\alpha)} = \frac{1}{2\pi} |(f_{\alpha}, g)|^2 \quad \text{for} \quad \alpha \in \mathbb{C}$$
(19)

where g is any vector of unit norm in H_q .

Because of (17), it follows that

$$p(\alpha, g) \ge 0, \qquad \int_{\alpha \in \mathbb{C}} d\mu(\alpha) p(\alpha, g) = 1$$
 (20)

hold for any vector g of unit norm in H_q .

3. PHASE DISTRIBUTION

To obtain the phase distribution we consider first the *phase operator* $P = (q^n + T^*T)^{-1/2}T$ and try to find the solution of the following eigenvalue equation:

$$Pf_{\beta} = \beta f_{\beta} \tag{21}$$

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We arrive at

$$f_{\beta} = \sum_{n=0}^{\infty} a_n \sqrt{[n]!} f_n$$

= $a_0 \sum_{n=0}^{\infty} \beta^n \sqrt{\frac{(q+[0])(q^2+[1])(q^3+[2])\dots(q^n+[n-1])}{[n]!}} f_n$

where $\beta = |\beta|e^{i\theta}$ is a complex number. These vectors are normalizable in a strict sense only for $|\beta| < 1$.

For details of the calculations see ref. 2.

Now, if we take $a_0 = 1$ and $|\beta| = 1$ we have

$$f_{\beta} = \sum_{n=0}^{\infty} e^{in\theta} \sqrt{\frac{(q+[0])(q^2+[1])(q^3+[2])\dots(q^n+[n-1])}{[n]!}} f_n \quad (22)$$

Henceforth, we shall denote this vector as

$$f_{\theta} = \sum_{n=0}^{\infty} e^{in\theta} \sqrt{\frac{(q+[0])(q^2+[1])(q^3+[2])\dots(q^n+[n-1])}{[n]!}} f_n \quad (23)$$

 $0 \le \theta \le 2\pi$, and call f_{θ} a phase vector in H_{q} .

The phase vectors f_{θ} are neither normalizable nor orthogonal. The completeness relation

$$I = \frac{1}{2\pi} \int_{X} \int_{0}^{2\pi} d\nu(x, \theta) |f_{\theta}\rangle \langle f_{\theta}|$$
(24)

where

$$d\nu(x,\,\theta) = d\mu(x)\,d\theta\tag{25}$$

may be proved as follows:

Here we consider the set X consisting of the points x = 0, 1, 2, ...,and $\mu(x)$ is the measure on X which equals

$$\mu_n \equiv \frac{[n]!}{(q + [0])(q^2 + [1]) \dots (q^n + [n - 1])}$$

at the point x = n; θ is the Lebesgue measure on the circle.

Define the operator

$$|f_{\theta}\rangle\langle f_{\theta}|: \quad H_q \to H_q$$
 (26)

by

$$\left|f_{\theta}\right\rangle \left\langle f_{\theta}\right| f = (f_{\theta}, f) f_{\theta} \tag{27}$$

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with $f(z) = \sum_{n=0}^{\infty} a_n z^n$. Now,

$$(f_{\theta}, f) = \sum_{n=0}^{\infty} [n]! \frac{e^{-in\theta}}{\sqrt{[n]!}} \sqrt{\frac{(q+[0])(q^2+[1])(q^3+[2])\dots(q^n+[n-1])}{[n]!}} a_n$$
$$= \sum_{n=0}^{\infty} e^{-in\theta} \sqrt{(q+[0])(q^2+[1])(q^3+[2])\dots(q^n+[n-1])} a_n \qquad (28)$$

Then,

$$(f_{\theta}, f)f_{\theta} = \sum_{n=0}^{\infty} \sum_{m=0}^{\infty} a_n e^{i(m-n)\theta} \sqrt{\frac{(q+[0])(q^2+[1])\dots(q^m+[m-1])}{[m]!}} \times \sqrt{(q+[0])(q^2+[1])\dots(q^n+[n-1])} f_m$$
(29)

Using

$$\int_{0}^{2\pi} d\theta \ e^{i(m-n)\theta} = 2\pi\delta_{mn} \tag{30}$$

we have

$$\begin{aligned} \frac{1}{2\pi} \int_{X} \int_{0}^{2\pi} d\nu(x,\,\theta) |f_{\theta}\rangle \langle f_{\theta}| f \\ &= \int_{X} d\mu(x) \sum_{n=0}^{\infty} \sum_{m=0}^{\infty} a_{n} f_{m} \sqrt{\frac{(q+[0])(q^{2}+[1])\dots(q^{m}+[m-1])}{[m]!}} \\ &\times \sqrt{(q+[0])(q^{2}+[1])\dots(q^{n}+[n-1])} \frac{1}{2\pi} \int_{0}^{2\pi} e^{i(m-n)\theta} d\theta \\ &= \sum_{n=0}^{\infty} a_{n} f_{n} \int_{X} \frac{(q+[0])(q^{2}+[1])\dots(q^{n}+[n-1])}{\sqrt{[n]!}} d\mu(x) \\ &= \sum_{n=0}^{\infty} a_{n} f_{n} \frac{(q+[0])(q^{2}+[1])\dots(q^{n}+[n-1])}{\sqrt{[n]!}} \\ &\times \frac{[n]!}{(q+[0])(q^{2}+[1])\dots(q^{n}+[n-1])} \end{aligned}$$

$$= \sum_{n=0}^{\infty} \sqrt{[n]!} a_n f_n$$
$$= f \tag{31}$$

Thus, (24) follows.

The *phase distribution* over the window $0 \le \theta \le 2\pi$ for any vector *f* is then defined by

$$P(\theta) = \frac{1}{2\pi} |(f_{\theta}, f)|^2$$
(32)

4. PHASE ESTIMATION

Once we have the POM information, we are ready to discuss the phase estimation problem. Without loss of generality, we assume that $0 \le \theta \le 2\pi$. The class of POMs we must optimize over in order to find the best phase estimate is taken to be

$$\{d\hat{\Pi}(\theta): 0 \le \theta \le 2\pi\}$$

where

$$d\hat{\Pi}(\theta) = d[\hat{\Pi}(\theta)]^{\dagger}$$
 and $I = \int_{0}^{2\pi} d\hat{\Pi}(\theta)$ (33)

The conditional probability density, given the phase operator

$$P = (q^n + T^*T)^{-1/2}T$$

for obtaining a phase value θ from this POM is

$$p(\theta, P) = \frac{(g, d \Pi(\theta)g)}{dv(x, \theta)}$$
 for $0 \le \theta \le 2\pi$, x an integer (34)

where g is a vector of unit norm in H_q .

We choose the POM $d\hat{\Pi}(\theta)$ and the input vector g to optimize our estimate of the phase shift P. For a given POM and the input vector, equation (34) supplies the PDF needed to perform a classical maximal likelihood estimation. The observed phase value θ is our estimate of P. In order for this estimate to be one of maximum likelihood, we restrict our attention to the POMs satisfying

$$P_{ML}(\theta) = \arg \max_{\theta} p(\theta, P) \quad \text{for } \theta \le 2\pi$$
 (35)

and optimize our estimate over $d \hat{\Pi}$ and g by maximizing the peak likelihood, minimizing $\delta \theta \equiv 1/p(\theta, P)$.

For the input vector

$$g = \sum_{n=0}^{\infty} (f_n, g) f_n$$

where

$$(f_n, g) = |(f_n, g)|e^{ik_n}, \qquad n = 0, 1, 2, \dots$$
 (36)

 $\delta\theta$ is minimized by the following POM:

$$d\hat{\Pi}(\theta) = \left| f_{\theta}^{g} \right\rangle \langle f_{\theta}^{g} \right| \frac{d\nu(x, \theta)}{2\pi}$$
(37)

where

$$d\nu(x, \theta) = d\mu(x) d\theta, \qquad 0 \le \theta \le 2\pi$$

as in (25) and

$$f_{\theta}^{g} \equiv \sum_{n=0}^{\infty} e^{in\theta + ik_{n}} \\ \times \sqrt{\frac{(q + [0])(q^{2} + [1])(q^{3} + [2])\dots(q^{n} + [n - 1])}{[n]!}} f_{n} \quad (38)$$

To calculate the reciprocal peak likelihood $\delta \theta$ with this optimum POM to estimate *P* we observe first

$$p(\theta, P) = \frac{(g, d\hat{\Pi}(\theta)g)}{d\nu(x, \theta)} = \frac{|(f_{\theta}^{g}, g)|^{2}}{2\pi}$$
$$= \frac{1}{2\pi} \left| \sum_{n=0}^{\infty} e^{-in\theta} \sqrt{\frac{(q+[0])(q^{2}+[1])\dots(q^{n}+[n-1])}{[n]!}} \left| (f_{n}, g) \right|^{2}$$
(39)

Hence a suitable peak likelihood $\delta\theta$ for maximum $p(\theta, P)$ can be [4] $\delta\theta = 2\pi |(f_{\theta}^{g}, g)|^{-2}$

$$= 2\pi \left| \sum_{n=0}^{\infty} \sqrt{\frac{(q+[0])(q^2+[1])\dots(q^n+[n-1])}{[n]!}} \left| (f_n,g) \right| \right|^{-2}$$
(40)

which is independent of the phases $\{k_n\}$. In fact, $p(\theta, P)$ is independent of the phases $\{k_n\}$.

As the peak likelihood $\delta\theta$ is independent of $\{k_n\}$, we can assume, without loss of generality, that the input vector $g = \sum_{n=0}^{\infty} (f_n, g) f_n$ has positive real coefficient (f_n, g) . Equation (38) then reduces to

$$f_{\theta}^{g} = f_{\theta} = \sum_{n=0}^{\infty} e^{in\theta} \sqrt{\frac{(q+[0])\dots(q^{n}+[n-1])}{[n]!}} f_{n}$$
(41)

for $0 \le \theta \le 2\pi$, which is the solution of the eigenvalue equation (21),

$$Pf_{\theta} = e^{i\theta}f_{\theta}$$

Now we consider the operator

$$U = \sum_{n=0}^{\infty} e^{-ik_n} |f_n\rangle \langle f_n|$$
(42)

Observe that

$$UU^* = U^*U = I$$

Thus, U is a unitary transformation.

Now, for an arbitrary input vector g the optimum POM from equation (37) is equivalent to performing the unitary transformation U followed by the POM

$$d\Pi(\theta) = |f_{\theta}\rangle \langle f_{\theta}| \frac{d\nu(x, \theta)}{2\pi}$$
(43)

where

$$d\nu(x, \theta) = d\mu(x) d\theta, \qquad 0 \le \theta \le 2\pi$$

as in (24) and (25), for

$$Uf_{\theta}^{g} = \sum_{n=0}^{\infty} e^{in\theta + ik_{n}} e^{-ik_{n}} \sqrt{\frac{(q + [0]) \dots (q^{n} + [n - 1])}{[n]!}} f_{n}$$
$$= \sum_{n=0}^{\infty} e^{in\theta} \sqrt{\frac{(q + [0]) \dots (q^{n} + [n - 1])}{[n]!}} f_{n}$$
$$= f_{\theta}$$
(44)

where f_{θ}^{g} is given by (38).

Shifting the input vector's phase by the phase operator P amounts to

$$(f_n, g) \to e^{in\theta_0}(f_n, g)$$
 for $n = 0, 1, 2, ...$ (45)

By rotating out the input phases k_n with the U transformation we get the transformed input as

$$e^{in\theta_0}(f_n, g) \xrightarrow{U} e^{in\theta_0} |(f_n, g)|$$
(46)

The effect of the POM on equation (43) on this transformed vector

$$g' = \sum_{n=0}^{\infty} e^{in\theta_0} |(f_n, g)| f_n$$
(47)

gives the classical phase with PDF

$$p(\theta, P) = \frac{(g', d \Pi(\theta)g')}{d\nu(x, \theta)}$$

= $\frac{1}{2\pi} (g', |f_{\theta}\rangle\langle f_{\theta}|g')$
= $\frac{|(f_{\theta}, g')|^{2}}{2\pi}$
= $\frac{1}{2\pi} \left| \sum_{n=0}^{\infty} e^{in(\theta_{0}-\theta)} \sqrt{\frac{(q+[0])\dots(q^{n}+[n-1])}{[n]!}} |(f_{n}, g)| \right|^{2}$ (48)

From the above equation it is clear that ML estimate obeys $P_{ML}(\theta) = \theta$.

Thus, the POM in equation (24) leads to the ML phase estimate for all vectors in H_q . Thus, to achieve our goal of jointly optimizing phase estimate performance over both the measurement and the input vector, it remains for us to minimize $\delta\theta$ from equation (40) by appropriate choice of input vector. Specifically, the coefficients { (f_n, g) } for the input vector must minimize the right side of equation (40) subject to the normalization constraint

$$\sum_{n=0}^{\infty} |(f_n, g)|^2 = 1$$
(49)

and the average number constraint

$$\sum_{n=0}^{\infty} [n] |(f_n, g)|^2 = N_0$$
(50)

where $N_0 = (g, T^*Tg)$.

Without loss of generality, we shall assume that (f_n, g) are positive real. Now, maximize

$$L(g, \lambda_1, \lambda_2) \equiv \frac{1}{2\pi} \left[\sum_{n=0}^{\infty} (f_n, g) \right]^2 + \lambda_1 \left[\sum_{n=0}^{\infty} (f_n, g)^2 - 1 \right] + \lambda_2 \left[\sum_{n=0}^{\infty} [n] (f_n, g)^2 - N_0 \right]$$
(51)

where λ_1 and λ_2 are Lagrange multipliers.

It is straightforward to show that

$$(f_n, g) = \frac{c}{k + [n]}$$
 for $n = 0, 1, 2, ...$ (52)

achieves the required stationary point for L, where c and k are positive constants depending on the Lagrange multipliers. For brevity we shall chose k = 1.

As we know that $[n] \ge n$ for q > 0, we have

$$\frac{c/(1+[n])}{1/n} \le \frac{c}{1/n+1}$$

Hence we see that

$$\lim_{n \to \infty} \frac{c/(1 + [n])}{1/n} \le c$$

Thus, the series $\sum_{n=0}^{\infty} [c/(1 + [n])]$ and $\sum_{n=0}^{\infty} (1/n)$ converge or diverge together. But $\sum_{n=0}^{\infty} (1/n)$ diverges. Hence, we must introduce a truncation parameter in equation (52). That is, we have

$$(f_n, g) = \frac{c}{1 + [n]}$$
 for $n = 0, 1, 2, ..., s$
= 0 for $n > s$ (53)

Now, we have

$$N_{0} = \sum_{n=0}^{s} [n] \cdot |(f_{n}, g(\alpha))|^{2}$$

= $\sum_{n=0}^{s} [n] \frac{c^{2}}{(1 + [n])^{2}}$
= $\sum_{n=0}^{s} \frac{c^{2}}{1 + [n]} - 1$ (54)

where we have used equations (50), (51), and (53) with the truncation point *s*. Then,

$$\delta\theta = 2\pi \left[\sum_{n=0}^{s} \sqrt{\frac{(q+[0])(q^2+[1])\dots(q^n+[n-1])}{[n]!}} (f_n,g) \right]^{-2}$$
$$= 2\pi c^2 \left[\sum_{n=0}^{s} \sqrt{\frac{(q+[0])(q^2+[1])\dots(q^n+[n-1])}{[n]!}} \frac{c^2}{1+[n]} \right]^{-2}$$

$$< 2\pi c^2 A \left[\sum_{n=0}^{s} \frac{c^2}{1+[n]} \right]^{-2} = \frac{2\pi c^2 A}{(N_0+1)^2} \approx \frac{2\pi c^2 A}{N_0^2}$$
(55)

for $N_0 \gg 1$. Here *A* is a constant.

5. CONCLUSION

We know [3] that ML phase estimation with the optimized state leads to $\delta\theta \sim 1/N_0^2$ for the reciprocal peak likelihood performance, where we are interested in the behavior at high average photon number, namely $N_0 \gg 1$. In this paper we show that in the deformed case $\delta\theta$ can be even less than $1/N_0^2$.

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